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Crystallographic polarity determination of $\gamma$-CuI. By A. S. Bhalla and E. W. White, Materials Research Laboratory, The Pennsylvania State University, University Park, Pa. 16802, U.S.A.
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#### Abstract

The polarity of single crystals of CuI having the sphalerite structure was determined from intensities measured on a double-crystal X-ray spectrometer, making use of anomalous scattering. If the direction from Cu to $I$ is defined as [111], the results confirm that the (111) plane contains $\mathrm{Cu}^{+}$ions while the (111) plane contains I-ions, and that the natural faces of the tetrahedral crystals are of the type (111).


Halides of $\mathrm{Cu}^{+}$have the sphalerite structure with space group $F \mathbf{4 3 m}$. The positive sense of the [111] direction is defined from $\mathrm{Cu}^{+}$to the halide ion, when $\mathrm{Cu}^{+}$is considered at position $(0,0,0)$ and the halide at position ( $\left(\frac{1}{4}, \frac{1}{4}, \frac{1}{4}\right)$. The opposite faces (111) and (II1) of this type of structure behave differently with regard to oxidation (Lavine, Rosenberg \& Gatos, 1958) and etching (Maringer, 1958; Schell, 1957) and show differences in the integrated X-ray intensities (Coster, Knol \& Prins, 1930) of reflection (Coster, Knol \& Prins, 1930; Geib \& Lark-Horowitz, 1932). The crystallographic polarity of CuCl and CuBr was established by Monier \& Kern (1955). In this note results on $\gamma$-CuI are reported and correlated with the above definition of the [111] direction.

Crystals of $\gamma-\mathrm{CuI}$ ( 1 mm diameter) were used for experimental measurements on a double crystal X-ray spectrometer (Bhalla \& White, 1970) having a finely collimated beam diameter smaller than the size of the crystal. Table 1 shows the calculated and measured intensity ratios from the two opposite faces, perpendicular to [111], for different orders of diffraction.

The observed ratios are very close to the calculated values. The X-ray conditions (counting time, current and voltage) were identical only while comparing a particular pair ( $h k l$ ) and ( $\bar{h} \bar{k} \bar{l}$ ).

Calculations were made considering $\mathrm{Cu}^{+}$at $(0,0,0)$ and $\mathrm{I}^{-}$at $\left(\frac{1}{4}, \frac{1}{4}, \frac{1}{4}\right)$ and using the formula (Warekois \& Metzger, 1959):

$$
\left|F_{n}\right|^{2}=16\left[f_{\mathrm{Cu}}+f_{\mathrm{I}}{ }^{i x x}\right]\left[f_{\mathrm{Cu}}^{*}+f_{\mathrm{I}}^{*} e^{-i x}\right]
$$

where $f_{\mathrm{Cu}}=f_{\mathrm{Cu}}^{0}+\Delta f_{\mathrm{Cu}}^{\prime}+i \Delta f_{\mathrm{Cu}}^{\prime \prime}$ etc. Asterisk symbols are the conjugate values of the complex quantities, and

$$
x=2 \pi \cdot \frac{3 n}{4}, n= \pm 1,2,3
$$

Values of $f^{\circ}, \Delta f^{\prime}$ and $\Delta f^{\prime \prime}$ were taken from International Tables for X-ray Crystallography (1967) and calculated for the appropriate order of diffraction.

The peaks, scanned from (111) and (111) faces, were found to have different full widths at half maximum. (Three sets of observations for first, second and third orders, in seconds of arc, are: 50 and 35,31 and 33, and 38 and 54 respectively.) These differences could be explained qualitatively from the prediction of the Darwin theory that, for a small angle of reflection from a perfect, nonabsorbing crystal, the peak half width at half maximum $\omega$, in radians, is related to the structure factor as given by (Compton \& Allison, 1935).

$$
\omega=2 \cdot 8 \delta F Z^{-1} \csc 2 \theta_{0}
$$

Table 1. Calculated and measured ratios of integrated intensities from opposite faces of $\gamma-\mathrm{CuI}$ single crystal for Cr K $\alpha$ radiation

Total counts (above background)

| $h k l$ | $h k l$ | $\bar{h} \bar{k} l$ |
| :---: | :---: | :---: |
| $(111)$ | 4332 | 2869 |
| $(222)$ | 1995 | 2029 |
| $(333)$ | 809 | 1723 |

$$
\frac{F F^{*}(h k l)}{F F^{*}(\bar{h} k \bar{l})}
$$

measured calculated

| 1.510 | 1.608 |
| :--- | :--- |
| 0.983 | 1.000 |

0.470
(In the parallel setting of the double crystal X-ray spectrometer, $Z=$ total number of electrons in the unit cell, $\theta_{0}$ is the Bragg angle and $\delta$ is the unit decrement of refractive index of the crystal.)

These X-ray diffraction intensity and peak width measurements show that the (111) plane of $\gamma$ - CuI contains $\mathrm{Cu}^{+}$ ions while the ( $\overline{1} \overline{1}$ ) plane contains $I^{-}$. Moreover, the studies showed that the natural faces of the tetrahedral crystals were of the ( $\overline{1} \overline{1} \overline{1}$ ) index. This study demonstrates the feasibility of working with single crystals as small as 1 mm in diameter for direct crystallographic polarity determinations using a double crystal X-ray spectrometer.

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$\mathrm{Ce}_{2} \mathrm{O}_{2} \mathrm{Sb}$ and $\mathrm{Ce}_{2} \mathrm{O}_{2} \mathrm{Bi}$ crystal structure.* By R. Benz, Los Alamos Scientific Laboratory, University of California, Los Alamos, New Mexico 87544, U.S.A.

$$
\text { (Received } 31 \text { August 1970) }
$$

$\mathrm{Ce}_{2} \mathrm{O}_{2} \mathrm{Sb}$ and $\mathrm{Ce}_{2} \mathrm{O}_{2} \mathrm{Bi}$ crystals were prepared with body-centered tetragonal $\left(\mathrm{U}_{2} \mathrm{~N}_{2} \mathrm{Sb}\right.$ type structure, space group $14 / \mathrm{mmm}$ ) unit-cell parameters

$$
\begin{array}{ll}
a_{0} \\
\mathrm{Ce}_{2} \mathrm{O}_{2} \mathrm{Sb} & 4.012 \pm 0.001 \AA \\
\mathrm{Ce}_{2} \mathrm{O}_{2} \mathrm{Bi} & 4.034 \pm 0.001
\end{array}
$$

$c_{0}$
$13.697 \pm 0.008 \AA$
$13.736 \pm 0.002$
$0.344+0.003$
$13.736 \pm 0.002$
$0.344 \pm 0.003$

Compounds with the chemical formula $\mathrm{Ce}_{2} \mathrm{O}_{2} \mathrm{X}, \mathrm{X}=\mathrm{Sb}$ and Bi , were prepared by reaction of weighed quantities of fused $\mathrm{Ce}_{2} \mathrm{O}_{3}+\mathrm{Ce}$ and of elemental X in sealed silica capsules for one month at $800^{\circ}$. X-ray diffraction powder patterns of the two products could be indexed on the basis of tetragonal symmetry. Unit-cell dimensions and densities calculated for two molecules per unit cell are listed in Table 1.

* Work done under the auspices of the U.S. Atomic Energy Commission.

The powder patterns are similar to that of the $\mathrm{U}_{2} \mathrm{~N}_{2} \mathrm{Sb}$ type previously reported for U and Th compounds (Benz \& Zachariasen, 1970) and to $\mathrm{Ce}_{2} \mathrm{O}_{2} \mathrm{Te}$ (Pardo \& Flahaut, 1962). Accordingly, the structure is

Space group $\mathrm{I} 4 / \mathrm{mmm}\left(D_{4 h}^{17}\right)$
(000) $\left(\frac{1}{2} \frac{1}{2}\right)+2 X$ on $(000)$

4 Ce in $\pm(00 u)$
$4 O$ in $\pm\left(\frac{1}{2} 0 \frac{1}{4}\right)$
where from observed intensities the parameter value was

Table 1. Lattice-parameter data and densities of $\mathrm{Ce}_{2} \mathrm{O}_{2} \mathrm{X}$ compounds

|  | $a_{0}$ | $b_{0}$ | $u$ | $v$ | Reference | $\begin{aligned} & \text { Density } \\ & \left(\mathrm{g} . \mathrm{cm}^{-3}\right) \end{aligned}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathrm{Ce}_{2} \mathrm{O}_{2} \mathrm{Sb}$ | $4.012 \pm 0.001 \AA$ | $13.697 \pm 0.008 \AA$ | $0.344 \pm 0.003$ | - | This work | $6.540 \pm 0.007$ |
| $\mathrm{Ce}_{2} \mathrm{O}_{2} \mathrm{Bi}$ | $4.034 \pm 0.001$ | $13.736 \pm 0.002$ | $0 \cdot 344 \pm 0.003$ | - | This work | $7.747 \pm 0.02$ |
| $\mathrm{Ce}_{2} \mathrm{O}_{2} \mathrm{Te}$ | $4.0705 \pm 0.0006$ | $12.904 \pm 0.002$ | $0.344 \pm 0.003$ | - | This work | $6.835 \pm 0.004$ |
|  | $4 \cdot 09$ | 12.92 |  |  | Pardo \& Flahaut (1962) | 6.76 |
| $\mathrm{Ce}_{2} \mathrm{O}_{2} \mathrm{~S}$ | $4.001 \pm 0.002$ | $6.844 \pm 0.006$ | $0.278 \pm 0.003$ | 0.63 | This work | $6.028 \pm 0.01$ |
|  | $4.01 \pm 0.01$ | $6.83 \pm 0.03$ | $0.29 \pm 0.02$ | $0 \cdot 64$ | Zachariasen (1949) | 5.99 |
|  | $4 \cdot 00$ | 6.87 | - | - | Flahaut \& Guittard (1955) | $6 \cdot 00$ |
|  | 4.004 | 6.872 | - |  | Picon \& Patrie (1956) | $6 \cdot 00$ |
|  | 4.008 | 6.886 | $0 \cdot 208$ | 0.863 | Flahaut, Guittard \& Patrie (1958) | $5 \cdot 98$ |
| $\mathrm{Ce}_{2} \mathrm{O}_{2} \mathrm{Se}$ | $4 \cdot 0152 \pm 0.0003$ | $7.0492 \pm 0.0009$ | $0.293 \pm 0.003$ | 0.63 | This work | $6 \cdot 603 \pm 0.002$ |
|  | $4.04 \pm 0.01$ | $7 \cdot 06 \pm 0 \cdot 01$ | - |  | Benacerraf, Guittard, Domange \& Flahaut (1959) | $6 \cdot 51$ |

